

## Effective Conductivity Concept for Modeling Conductor Surface Roughness

G. Gold, K. Helmreich

Institute of Microwaves and Photonics

Friedrich-Alexander-University Erlangen-Nuremberg

[gerald.gold@fau.de](mailto:gerald.gold@fau.de)

[klaus.helmreich@fau.de](mailto:klaus.helmreich@fau.de)



FRIEDRICH-ALEXANDER  
UNIVERSITÄT  
ERLANGEN-NÜRNBERG

TECHNISCHE FAKULTÄT

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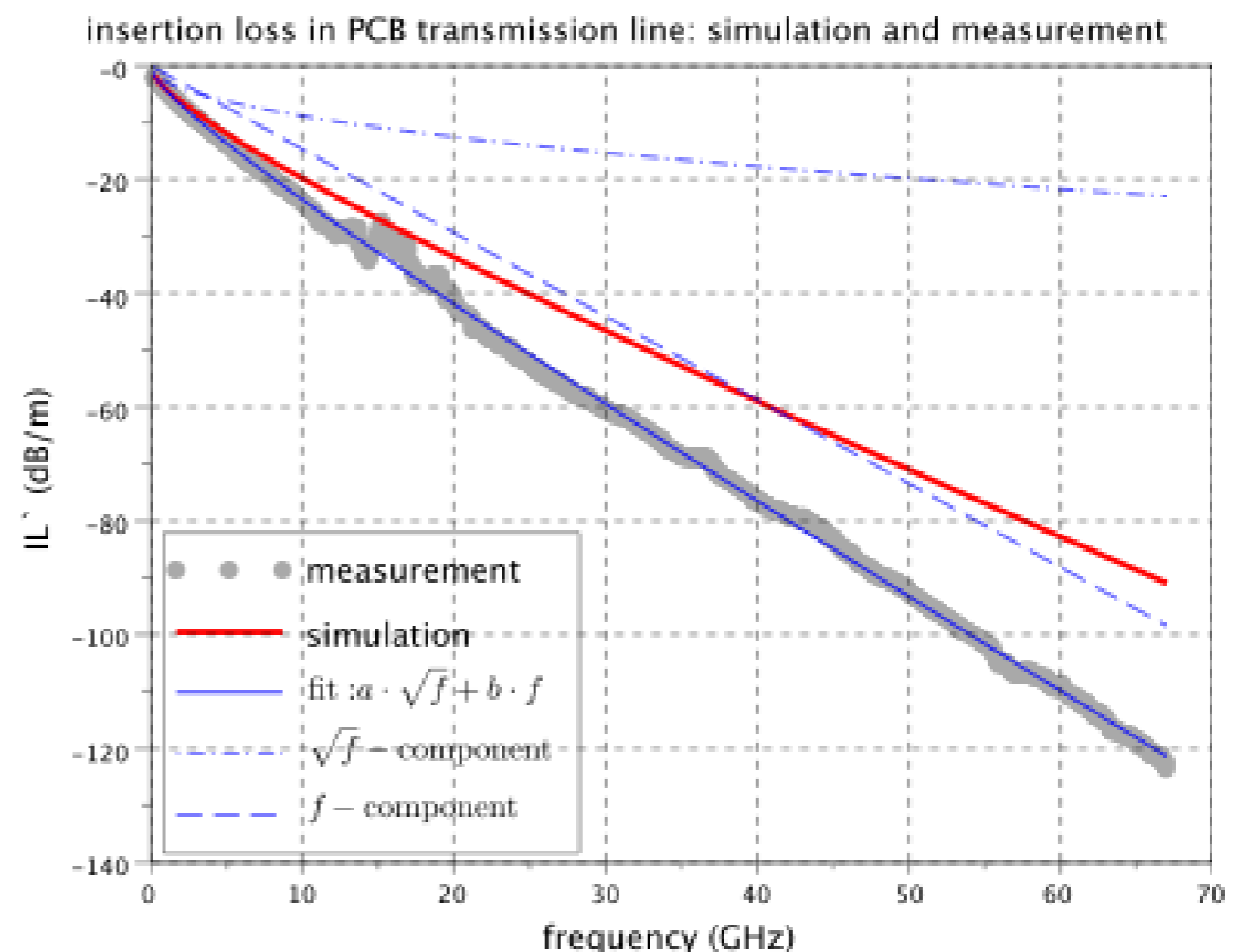
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# Introduction

- ◆ Accurate models of conductor behavior are essential for predictive simulations: antennas, resonators, connectors, filters, transmission lines
- ◆ Today, PCB transmission lines are operated up to the two digit Gigahertz range
- ◆ **Transmission line loss:**
- ◆ Usually dominated by dielectric loss  $\propto f$
- ◆ Theory predicts  $a\sqrt{f}$  - dependence due to skin effect for conductor loss
- ◆ **Skin effect at frequencies > 1GHz:**
- ◆ Skin depth decreases to the order of surface roughness:
  - conductor surface can no longer be regarded as ideally smooth
  - $\sqrt{f}$ - dependence is no longer valid



# Existing Models for Surface Roughness

## ◆ Phenomenological models:

◆ Correction factor  $K$  adapted to measurement

◆ Function of RMS roughness  $R_q$

• Hammerstad & Jensen  $K_{HJ} = 1 + \frac{2}{\pi} \tan^{-1} \left( 1.4 \frac{R_q}{\delta} \right)$

• Groiss  $K_G = 1 + \exp \left( - \left( \frac{\delta}{2R_q} \right)^{1.6} \right)$

• Fail at high frequencies resp. high values of  $R_q$

## ◆ Physical models:

◆ Huray's „snowball“-modell:

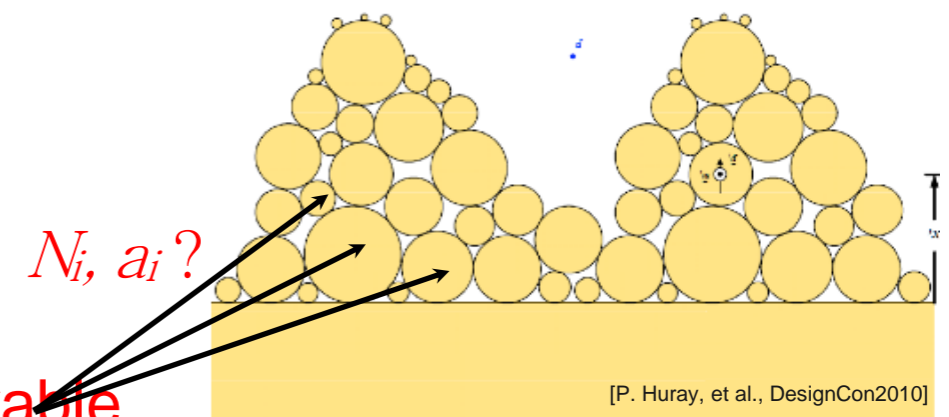
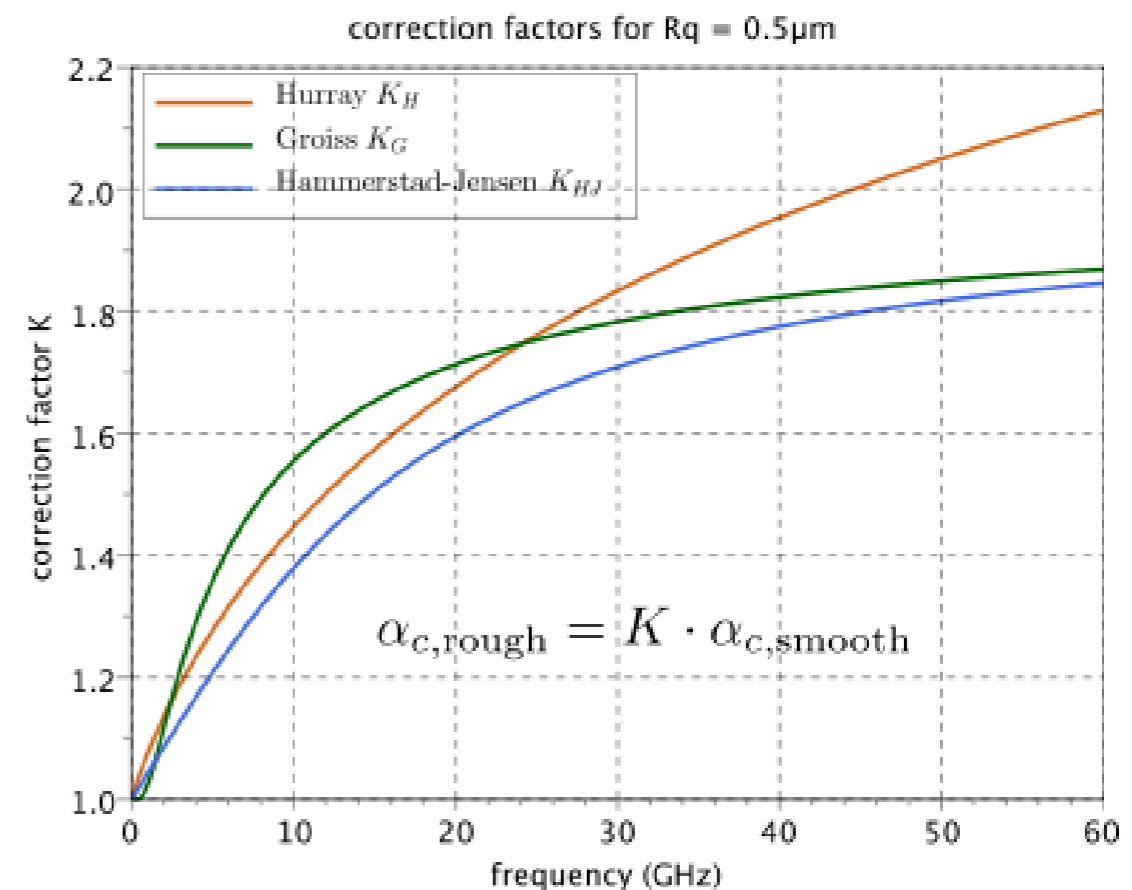
analytical calculation of additional power loss due to copper „snowballs“

$$K_H = 1 + \frac{3}{2} \sum_{i=1}^j \left( \frac{N_i}{A_{hex} \cdot 4\pi a_i^2} \right) \cdot \left( 1 + \frac{\delta}{a_i} + \frac{\delta^2}{a_i^2} \right)^{-1}$$

## ◆ Others:

◆ Fractal surface, „brute force“ 3D simulation, ...

◆ major drawbacks: many parameters, often hardly observable





# Inconsistencies of Common Modeling Approaches

## Current 'indirection'

- ◆ The displacement amplitude of a conduction electron is only  $\approx 10^{-12}\text{m}$  ( $P = 10\text{mW}$  at  $1\text{GHz}$ )
- ◆ There is no 'current indirection' in a rough surface in a sense that the current path follows any 'surface contour'

## Model roughness as piecewise smooth facets

- ◆ Surface profiles cannot be modeled by ideally smooth facets because then the solver **assumes the skin effect for an infinite plane** on their surfaces
- ◆ The plane skin effect is only valid if feature sizes  $\gg \lambda$

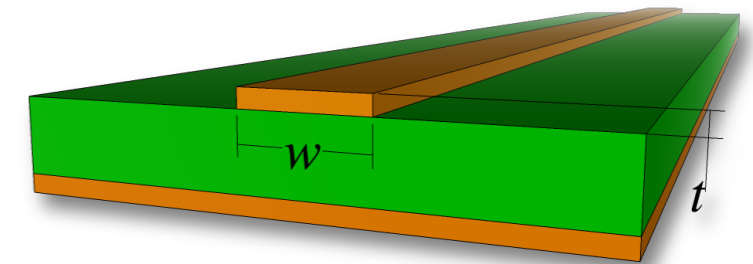
## Model microscopic features

- ◆ It is not necessary to simulate surface roughness on a **microscopic** level because it is **averaged over the order of a wavelength** anyway

# Dimensions Consideration

## Relevant dimensions for $f \approx 1-100\text{GHz}$ :

- ◆ Wavelength:  $\lambda \approx 2 - 200\text{mm}$
- 1. Trace width:  $w \approx 100\mu\text{m}$
- ◆ Trace thickness:  $t \approx 18\mu\text{m}$
- ◆ Skin depth (in Cu):  $\delta \approx 0.2 - 2\mu\text{m}$

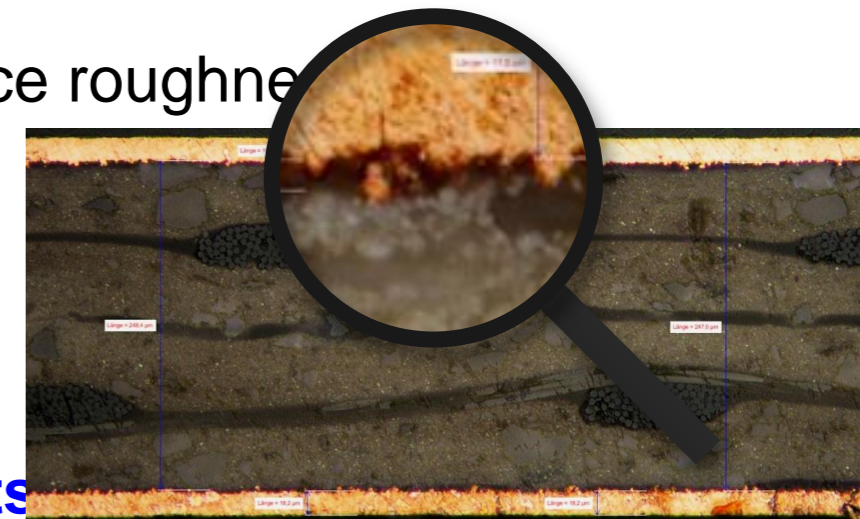


## Situation on PCB in the operating frequency range:

- Wave length  $\lambda \gg$  conductor dimensions  $w, t$
- Conductor dimensions  $w, t \gg$  skin depth  $\delta$  Skin depth  $\delta \lesssim$  surface roughness

## Conclusion:

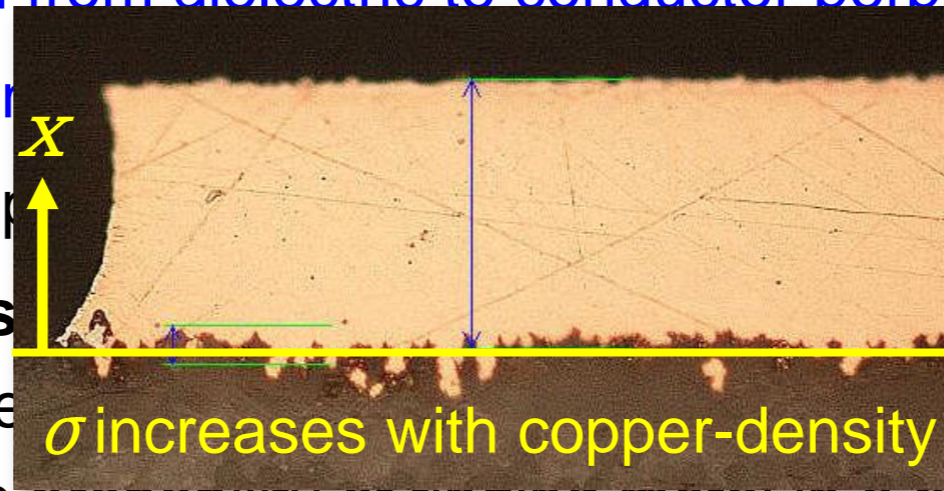
- ◆ As  $w, t \gg R_q$ , the conductor surface basically is “plane”
- ◆ But there is roughness on a microscopic scale ( $\ll \lambda, w, t$ )
- ◆ Propagating wave does not “see” individual peaks and pits however with no abrupt border between dielectric and conductor





# Modeling Approach

- ◆ No abrupt border between dielectric and conductor
- ◆ Not necessary to model **microscopic** peaks
- ◆ Rather model the **transition from dielectric to conductor perpendicular to the surface**
- ◆ Maintain **translation invariance**
- ◆ Appropriate **macroscopic** parameters
- ◆ **model surface roughness**
  - ▶  $\sigma$  is a function of distance  $x$
  - ▶  $\sigma(x)$  is proportional to the probability of finding metal in a plane parallel to the surface
  - ▶  $\sigma(x)$  increases from virtually zero in the dielectric to bulk metal conductivity



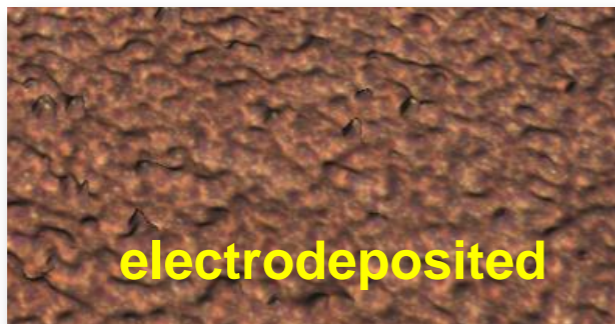
transition: **conductivity  $\sigma$**   
**perpendicular to the surface:**

# Surface Roughness Characterization

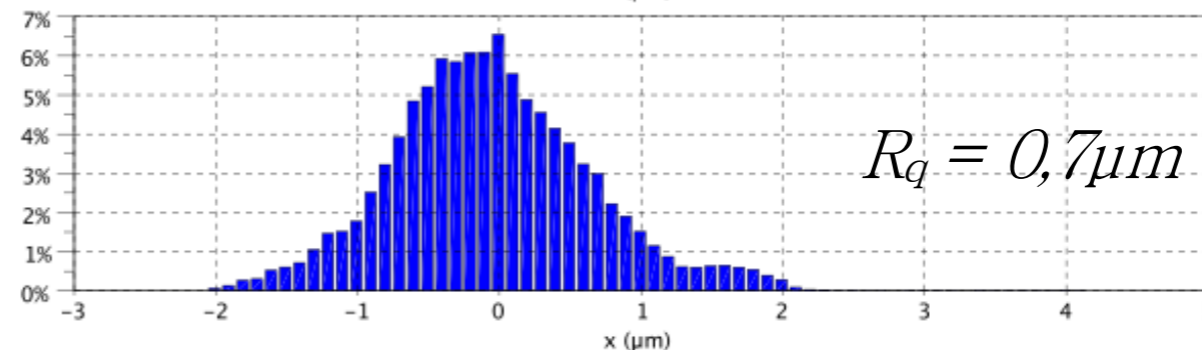
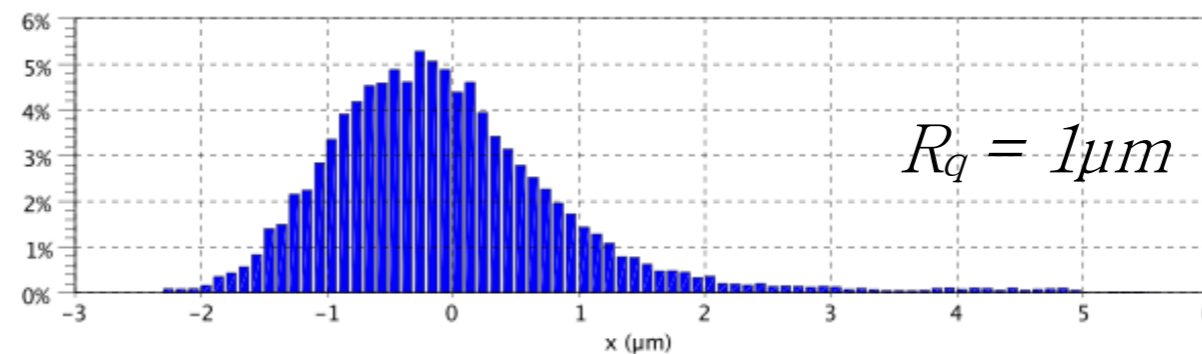
- ◆ Assume normally distributed surface profile
- ◆  $\sigma(x)$  corresponds to the cumulative distribution function (CDF) of surface profile:

$$\sigma(x) = \sigma_{bulk} \cdot \int_{-\infty}^x PDF(x) du = \sigma_{bulk} \cdot \int_{-\infty}^x \exp\left(-\frac{u^2}{2R_q^2}\right) du$$

- ◆ **Surface characterisation:**
  - ◆ Optical scanning system: **Confirms normal distribution**
- ⇒ **One single model parameter !**



measured PDF





# Conductivity Gradient Model Derivation

Skin effect in rough surfaces is deduced from Maxwell's equations with time harmonic fields and location dependent conductivity  $\sigma(r)$ :

Using Maxwell-Ampere's Law:  $\underline{\nabla} \times \left( \frac{\underline{B}}{\mu_r} \right) = j\omega \frac{1}{c^2} (\epsilon_r \underline{E}) + \mu_0 \underline{J}$ , inserting Ohm's Law  $\underline{J} = \sigma \underline{E}$

for  $\mu_r = 1$  yields: 
$$\underline{\nabla} \times \underline{B} = \left( j\omega \frac{\epsilon_r}{c^2} + \mu_0 \sigma \right) \underline{E}$$

Even at  $f = 100\text{GHz}$ ,  $\omega \epsilon_r / c^2 \ll \mu_0 \sigma$ , if  $\sigma \gg 5.6\text{S/m}$ , so displacement current density can be neglected for  $\sigma$  down to  $\approx 1\text{ppm}$  of copper conductivity:

$$\underline{\nabla} \times \underline{B} = \mu_0 \sigma \underline{E}$$

Taking the curl,  $\underline{\nabla} \times$ :


$$\underline{\nabla} \times \underline{\nabla} \times \underline{B} = \mu_0 \underline{\nabla} \times (\sigma \underline{E})$$

$$\underline{\nabla} \times \underline{\nabla} \times \underline{A} = \underline{\nabla}(\underline{\nabla} \cdot \underline{A}) - \Delta \underline{A} \quad / \quad \underline{\nabla} \times (\varphi \underline{A}) = \varphi \underline{\nabla} \times \underline{A} - \underline{A} \times \underline{\nabla} \varphi$$

we get

$$-\Delta \underline{B} = \mu_0 \sigma \underline{\nabla} \times \underline{E} - \mu_0 \underline{E} \times \underline{\nabla} \sigma$$

which, with Faraday's Law  $\underline{\nabla} \times \underline{E} = -j\omega \underline{B}$   $\left| \begin{array}{l} \underline{E} = \frac{1}{\mu_0 \sigma} \underline{\nabla} \times \underline{B} \end{array} \right.$

results in:  
$$\Delta \underline{B} - j\omega \mu_0 \sigma \underline{B} + \frac{\underline{\nabla} \sigma}{\sigma} \times (\underline{\nabla} \times \underline{B}) = 0$$

# Skin Effect in Rough Surfaces

- Assume gradual transition to bulk conductivity perpendicular to the mean surface
  - Focus on **one-dimensional** problem

## ◆ The gradient model

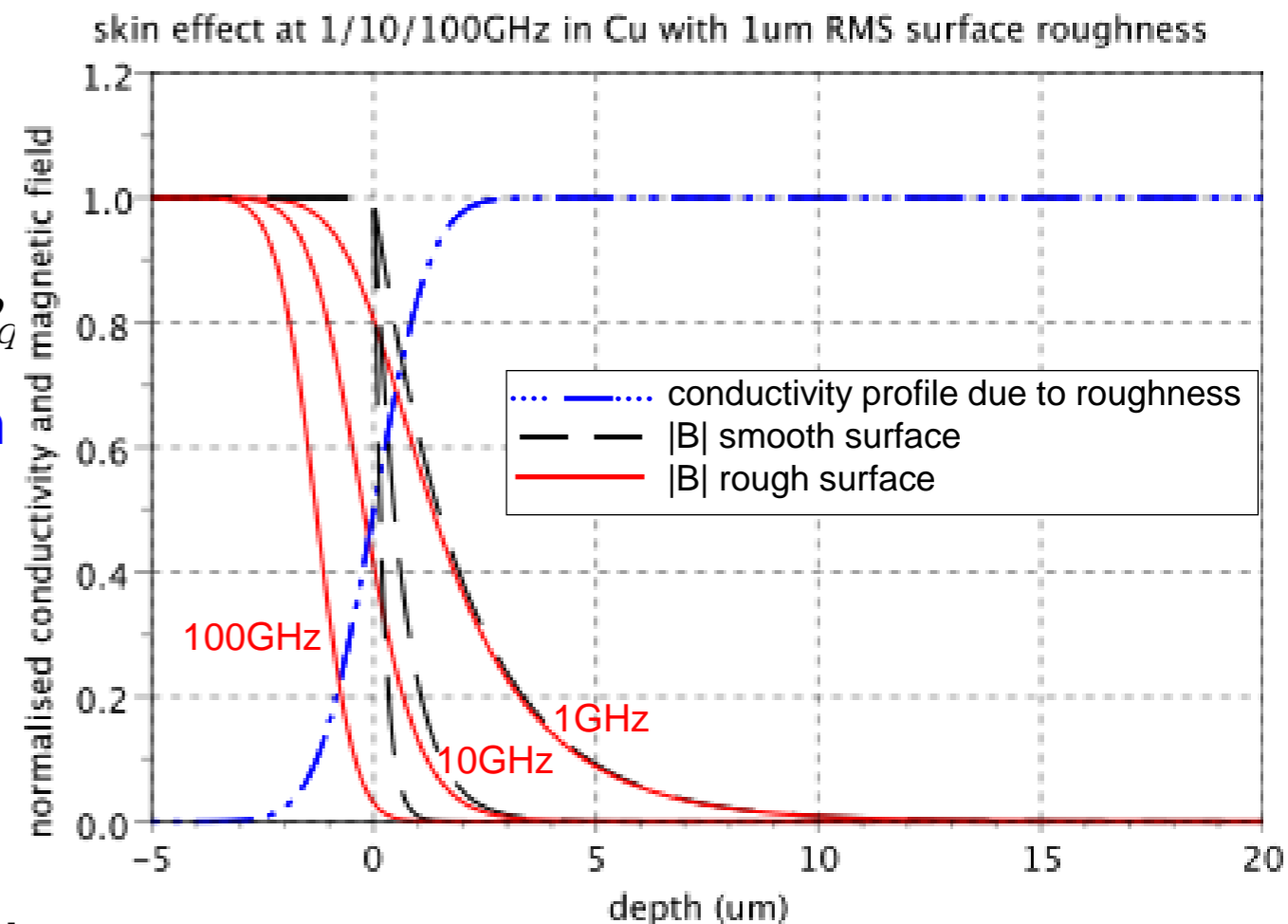
- Not only delivers a correction factor
- ◆ Describes skin-effect in rough surfaces as a whole by **one parameter**: RMS-roughness  $R_q$
- Calculates profiles of **magnetic field strength** and loss power density for a given roughness distribution

## ◆ Skin effect in ideally smooth surface:

- ◆ Magnetic field abruptly starts to decline

## ◆ Gradient model for skin effect in rough surfaces:

- ◆ Magnetic field smoothly decreases as it enters the range of surface roughness

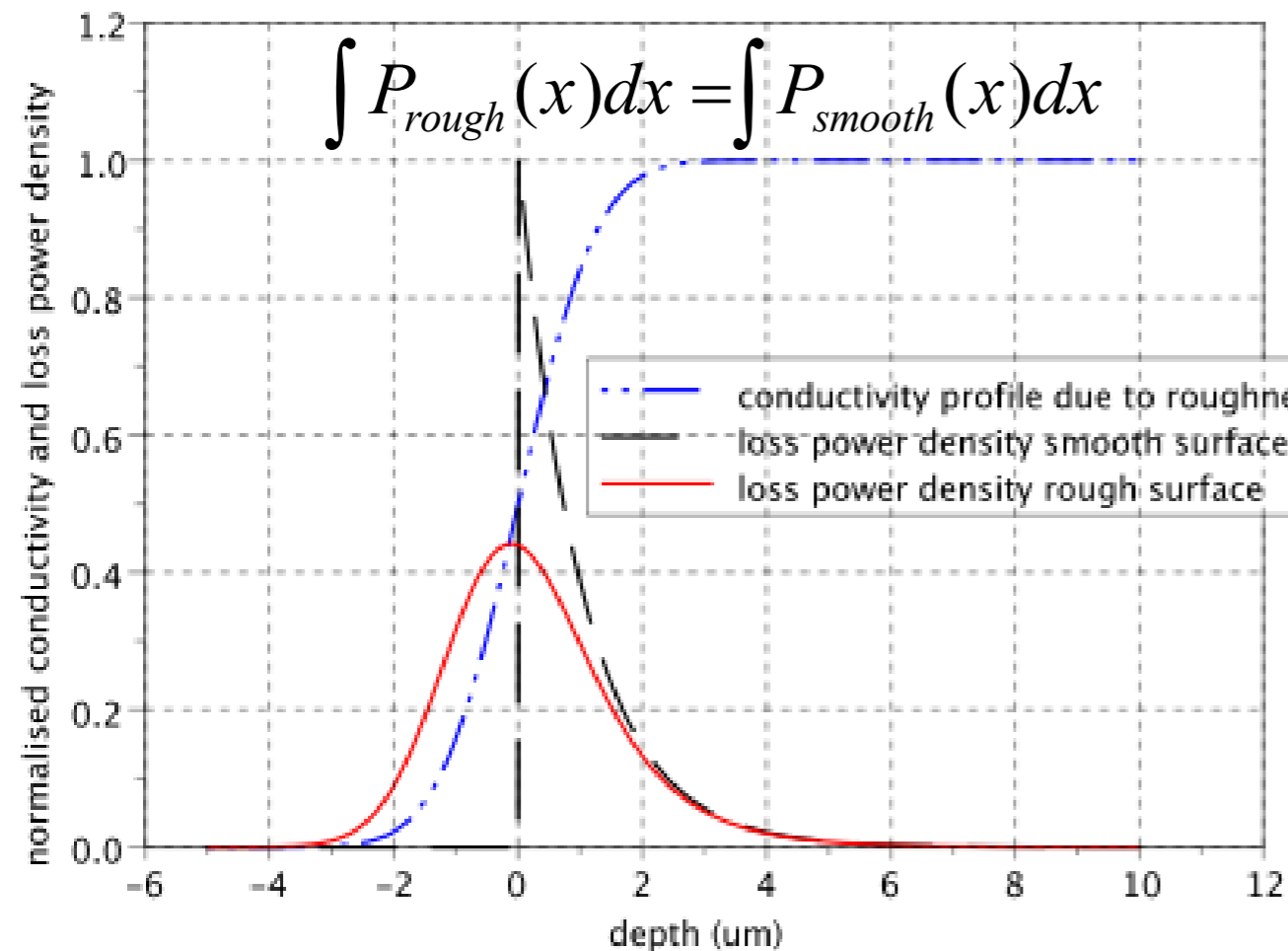




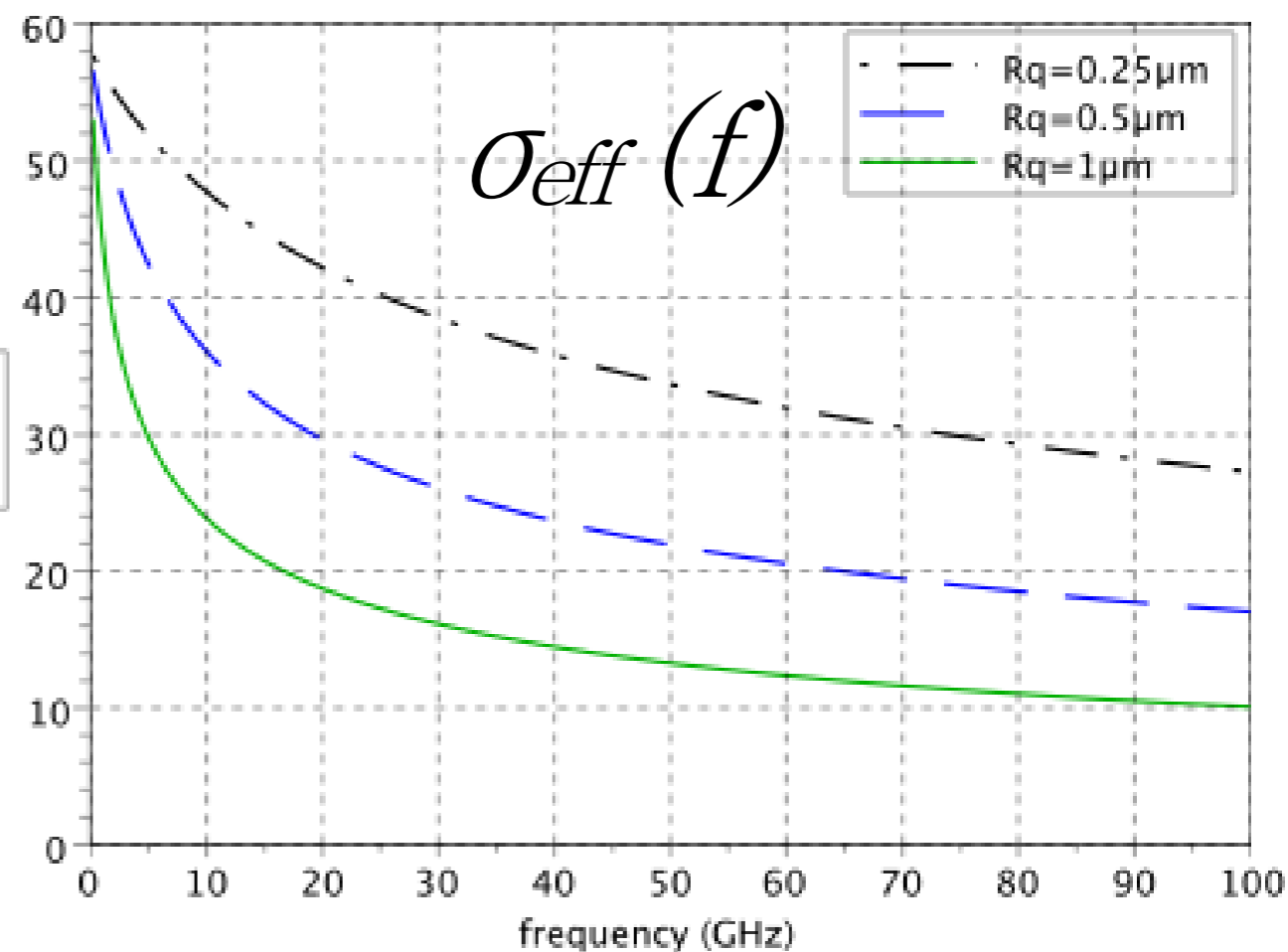
# Concept of Effective Conductivity

- ◆ Comparison of loss power densities: Gradient Model vs. conventional skin effect
- ◆ **Effective conductivity** is defined as the conductivity of a material with ideally smooth surface that would cause the same loss as the rough surface

skin effect at 1GHz in Cu with 1um RMS surface roughness



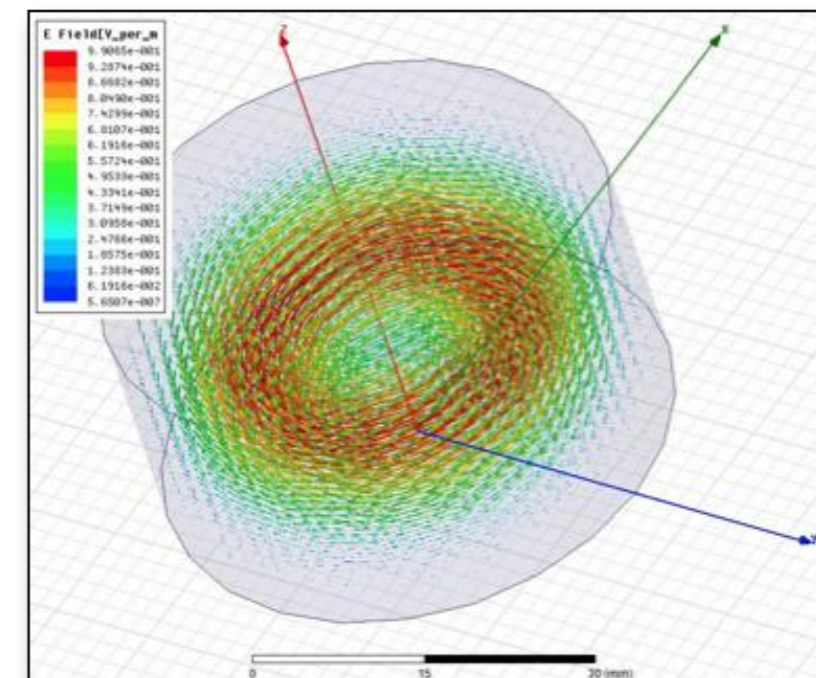
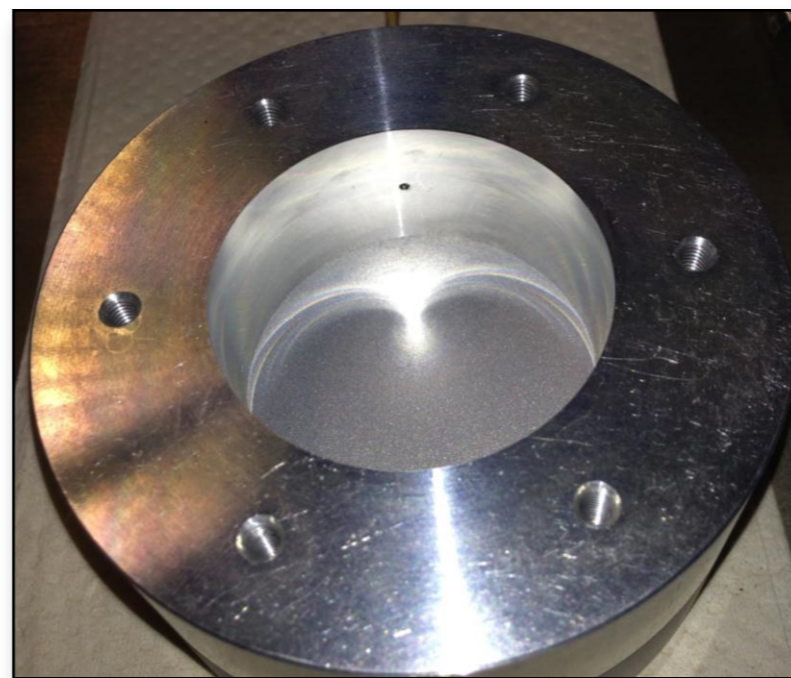
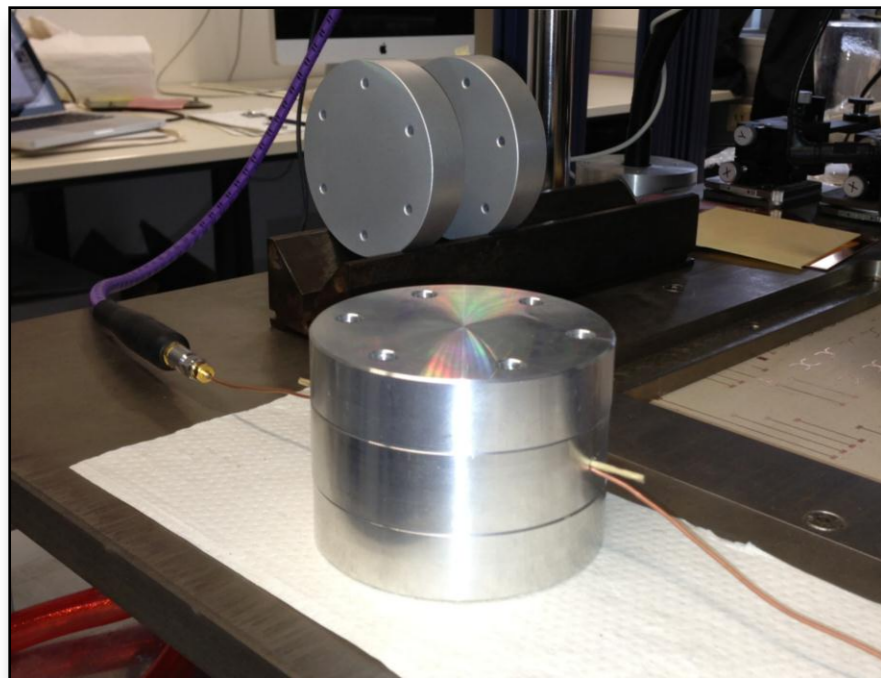
frequency dependent effective conductivity





# Direct Measurement of the Effective Conductivity

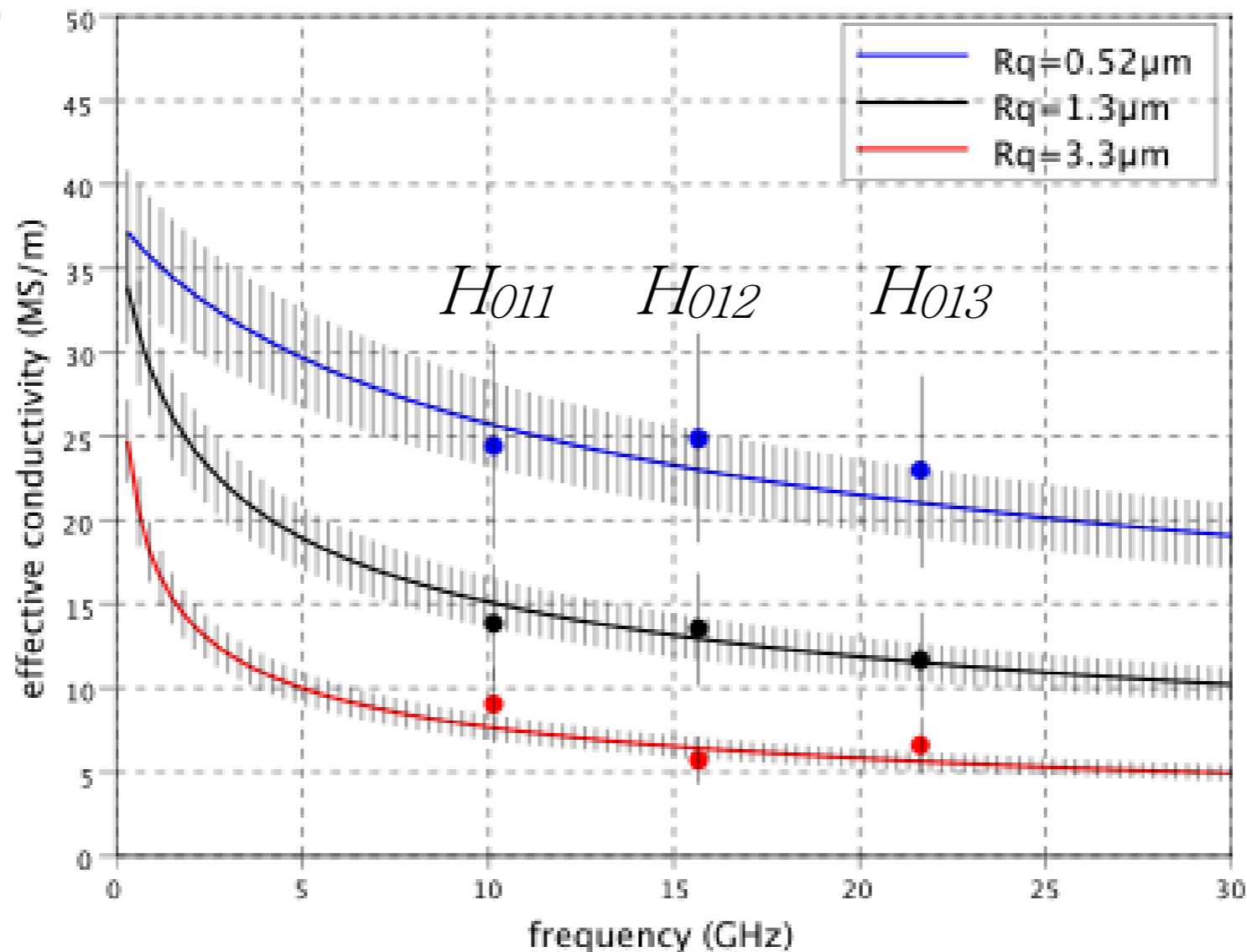
- ◆ Cylindrical cavity resonator operated at  $\approx 10, 15, 21\text{GHz}$  ( $H_{01x}$  modes)
- ◆ Electrical field has circular component  $E_\phi$  only
- ◆ Exchangeable lids of different surface roughness  $R_q$
- ◆ Influence of surface roughness measurable by quality factor  $Q$
- ◆ Treating rough surfaces as if they were ideally smooth yields  $\sigma_{eff}$





# Direct Measurement of the Effective Conductivity

- ◆ RMS-roughness  $R_q$  of lids was obtained by optical scanning system
- ◆ Responses for  $\sigma_{eff}$  predicted by Gradient Model from  $R_q$  agree with measured  $\sigma_{eff}$ :



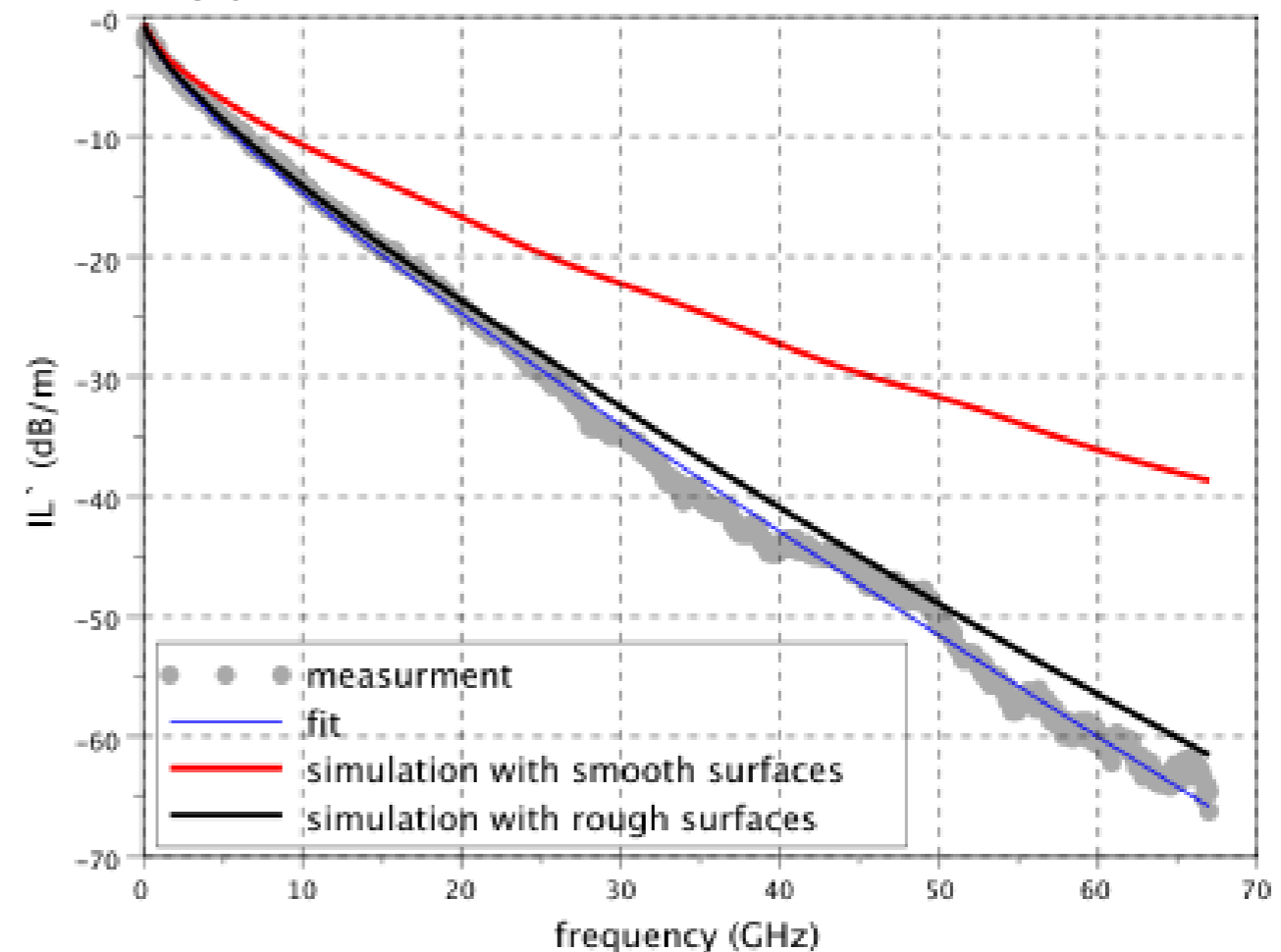
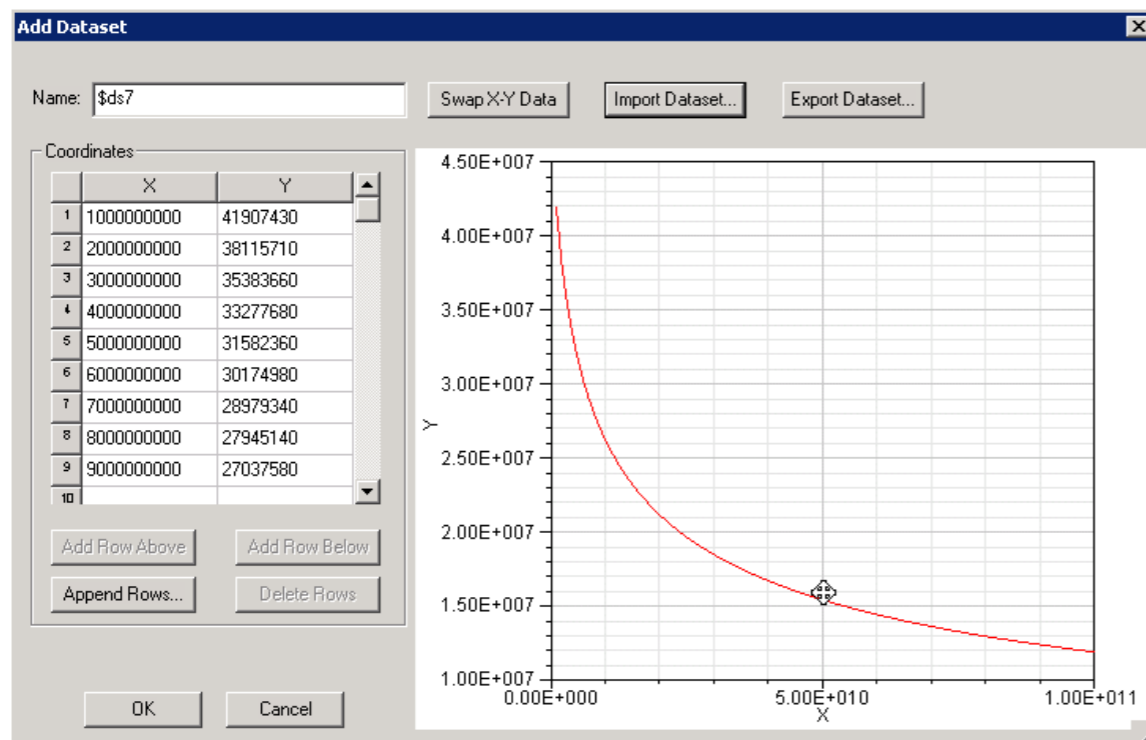
# Application in Field Solvers as Impedance Boundary Condition

Advantages of impedance boundary conditions:

- ◆ Not necessary to mesh inside the conductor
- ◆ No increase of computation time
- ◆ Impedance boundary condition with in the conductor to its surface

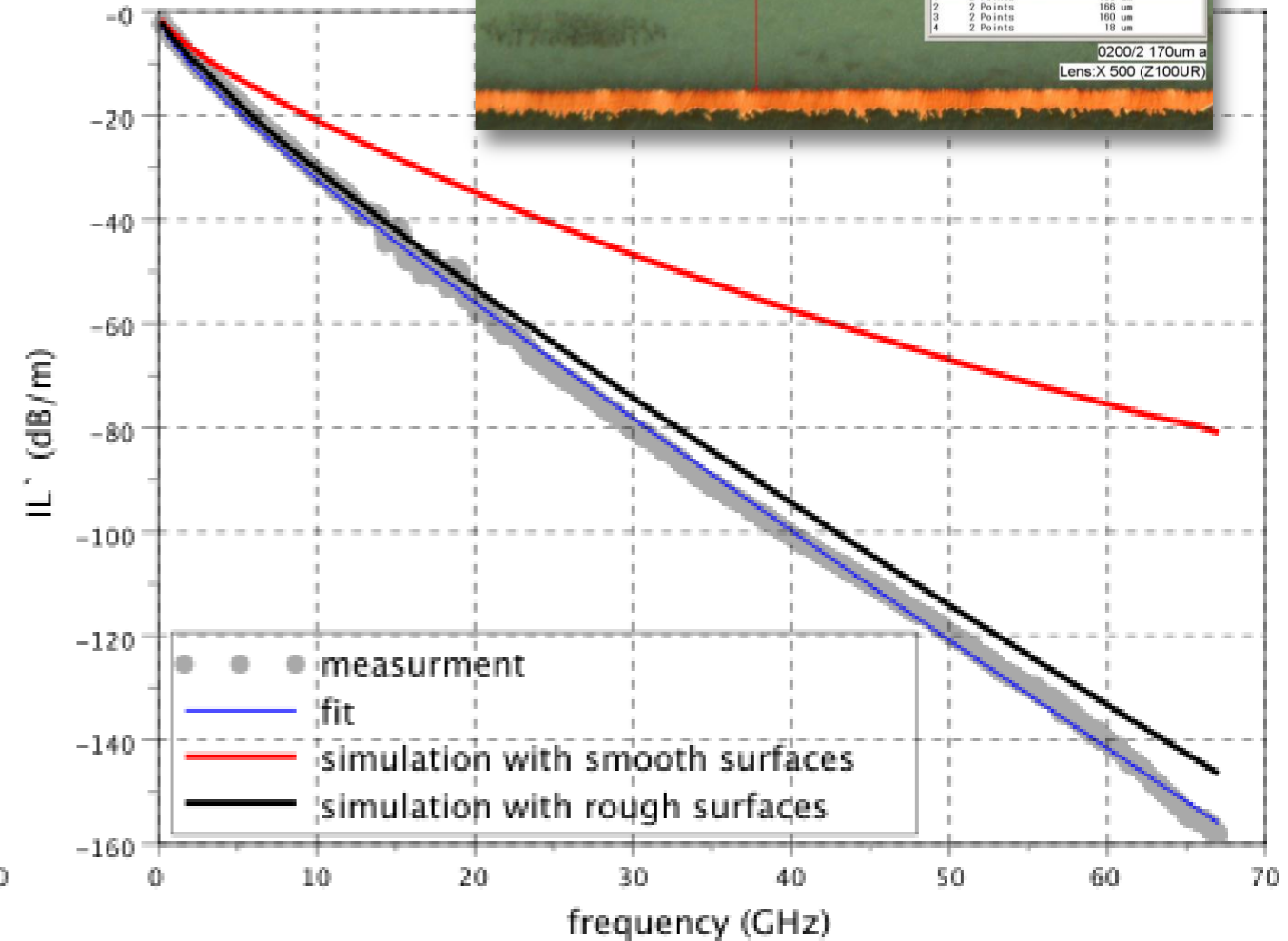
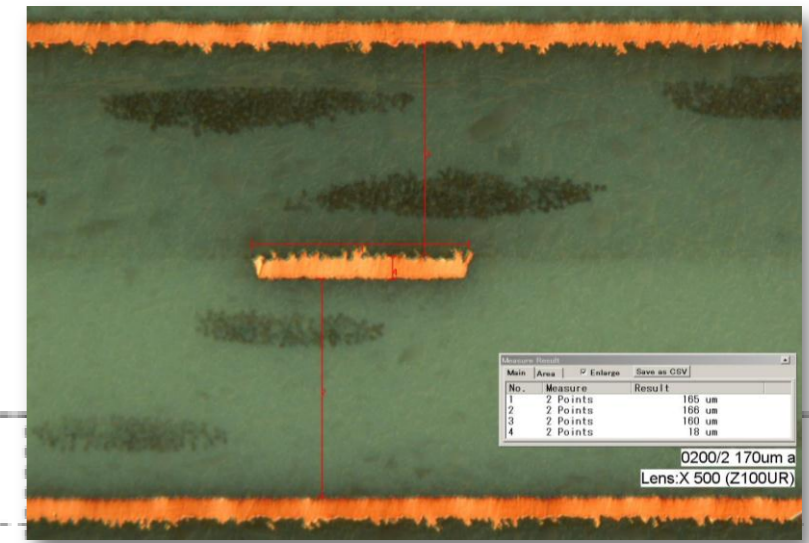
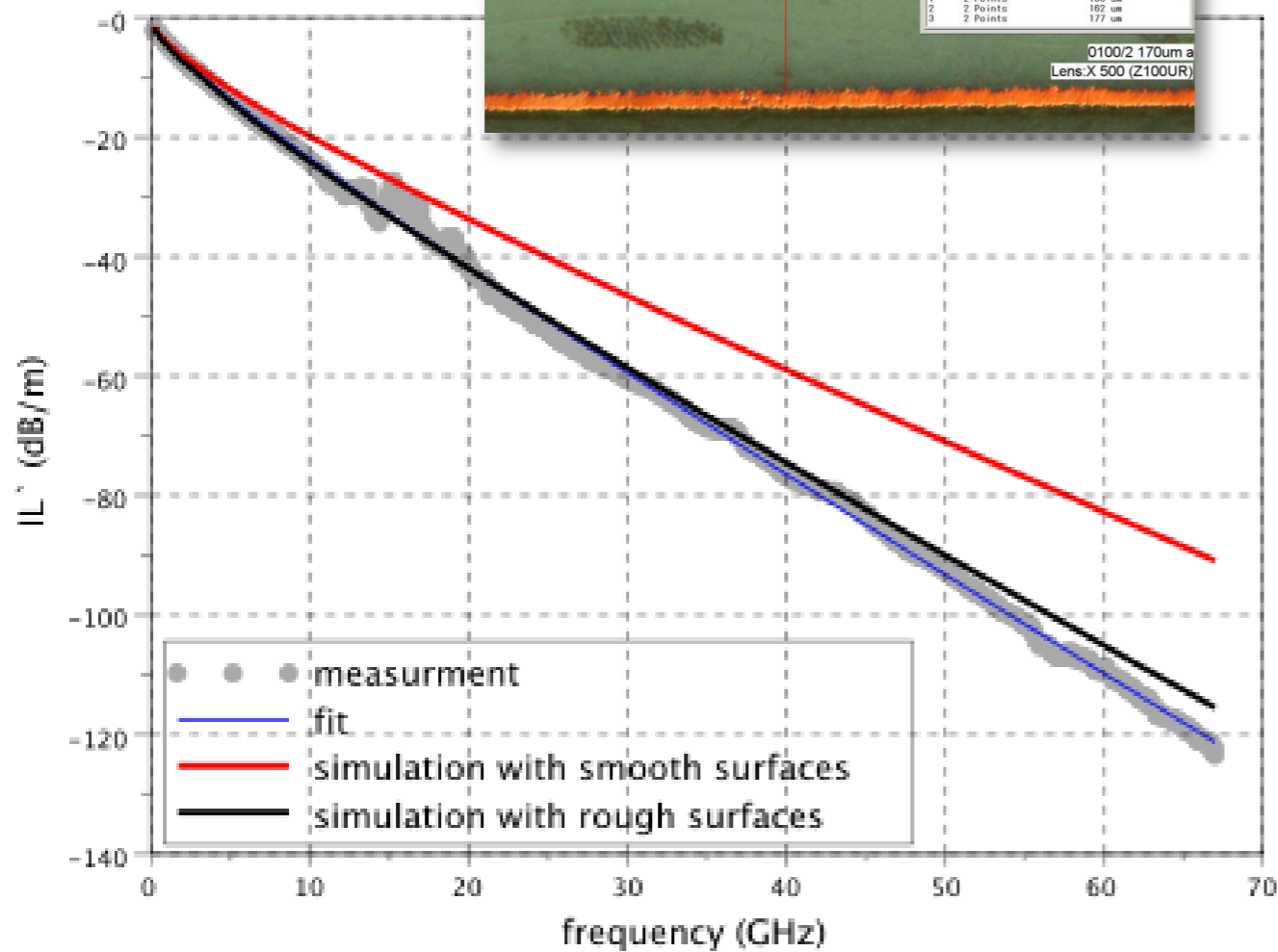
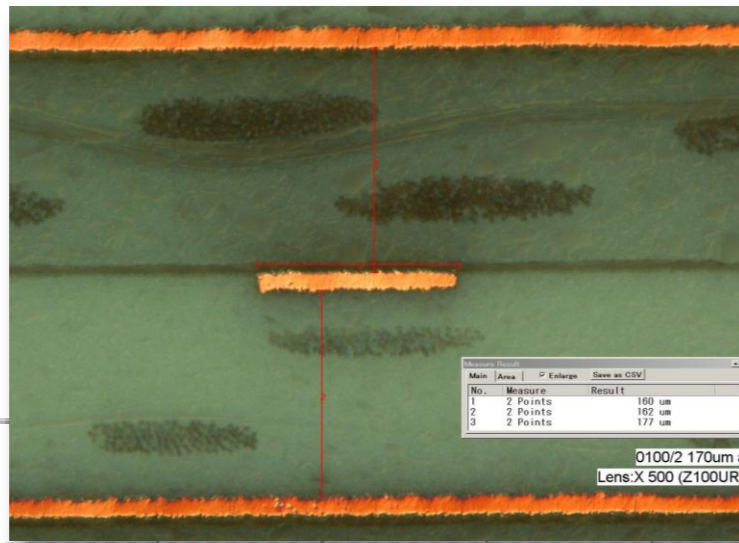
➔ Modification:  $\sigma \rightarrow \sigma_{eff}(f)$ ,  $\delta \rightarrow \delta(\sigma_{eff})$  !

$$Z_s = \frac{1+i}{\sigma\delta} \text{ projects field phenomena}$$





# Application in Field Solvers as Impedance Boundary Condition



# Conclusion

- ◆ Surface roughness is modeled as a **conductivity gradient**  $\sigma(x)$ 
  - ▶  $\sigma(x)$  is **proportional to the CDF** of the surface roughness profile
  - ▶ **Single parameter** model:  $R_q$
  - ▶  $R_q$  **measurable** with optical scanning system or microsection
  - ▶ **Datasheets** often provide values for  $R_q$
- ◆ Frequency dependent **effective conductivity**  $\sigma_{eff}$ 
  - ▶ Derived from comparison to loss power density of smooth surface
  - ▶ Surface impedance **as boundary condition**
  - ▶ **Easily applicable** with commercial field solvers: import once to library
  - ▶ No increase of computation time